

# Application of Laser Induced Electron Impact Ionization to the Deposition Chemistry in the Hot-Wire Chemical Vapor Deposition Process with $\text{SiH}_4$ - $\text{NH}_3$ Gas Mixtures

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The application of a laser-induced electron impact (LIEI) ionization source in studying the gas-phase chemistry of the  $\text{SiH}_4/\text{NH}_3$  hot-wire chemical vapor deposition (HWCVD) system has been investigated. The LIEI source is achieved by directing an unfocused laser beam containing both 118 nm (10.5 eV) vacuum ultraviolet (VUV) and 355 nm UV radiations to the repeller plate in a time-of-flight mass spectrometer. Comparison of the LIEI source with the conventional 118 nm VUV single-photon ionization (SPI) method has demonstrated that the intensities of the chemical species with ionization potentials (IP) above 10.5 eV, e.g.,  $\text{H}_2$ ,  $\text{N}_2$  and He, have been significantly enhanced with the incorporation of the LIEI source. It is found that the SPI source due to the 118 nm VUV light coexists in the LIEI source. This allows simultaneous observations of parent ions with enhanced intensity from VUV SPI and their "fingerprint" fragmentation ions from LIEI. It is, therefore, an effective tool to diagnose the gas-phase chemical species involved with both  $\text{NH}_3$  and  $\text{SiH}_4$  in the HWCVD reactor. In using the LIEI source to  $\text{SiH}_4$ ,  $\text{NH}_3$  and their mixtures, it has been shown that the  $\text{NH}_3$  decomposition is suppressed with the addition of  $\text{SiH}_4$  molecules. Examination of the  $\text{NH}_3$  decomposition percentage and the time to reach the  $\text{N}_2$  and  $\text{H}_2$  steady-state intensities for various  $\text{NH}_3/\text{SiH}_4$  mixtures suggests that the extent of the suppression is enhanced with more  $\text{SiH}_4$  content in the mixture. With increasing filament temperatures, the negative effect of  $\text{SiH}_4$  becomes less important. (J Am Soc Mass Spectrom 2007, 18, 1950–1958) © 2007 American Society for Mass Spectrometry

Silicon nitride ( $\text{SiN}_x$ ) thin films, deposited by hot-wire chemical vapor deposition (HWCVD) methods at low substrate temperatures, have attracted increasing research interest owing to their wide range of applications, such as anti-reflection coatings, passivation layers, diffusion barriers, and gate insulators, in semiconductor and electronics industry [1, 2]. The two simple hydrides,  $\text{SiH}_4$  and  $\text{NH}_3$ , are the most widely used source gases for  $\text{SiN}_x$  film formation using HWCVD. The deposition chemistry from  $\text{SiH}_4$  has been extensively studied in light of its use as the source gas for silicon thin-film deposition [3, 4]. At the high filament temperatures typical of Si film deposition, Si and H atoms are the primary species desorbed from the filament [5]. The gas-phase reactions of the two primary radicals with ambient silane molecules produce the final mix of depositing precursors such as  $\text{SiH}_3$  and  $\text{Si}_2\text{H}_{2y}$  ( $y = 1-3$ ) [6–9]. Ammonia decomposition chemistry, using tungsten metal in particular, has been studied since the early days of heterogeneous catalysis [10]. In studying the ammonia plasma processing of

metal and polymer substrates, Butoi et al. demonstrated that  $\text{NH}_2$  radicals are generated at the substrate surfaces [11]. Recently, Umemoto and coworkers studied the decomposition of  $\text{NH}_3$  on a hot tungsten filament [12]. They identified H atom and  $\text{NH}_2$  radical as primary products of ammonia decomposition. Two important stable products,  $\text{H}_2$  and  $\text{N}_2$ , were observed using mass spectrometry.

Although knowledge of the HWCVD processes with two components has been obtained, details of gas-phase chemistry involved with  $\text{SiH}_4$ - $\text{NH}_3$  gas mixtures have been much less studied and, consequently, a full understanding has not yet been achieved. Using deuterated silane ( $\text{SiD}_4$ ) as a source gas, Stannowski et al. [13] has shown that  $\text{SiH}_4$  is dissociated effectively at the filaments, whereas  $\text{NH}_3$  is scarcely dissociated at the filaments but rather in the gas phase by secondary reactions. Later, Umemoto et al. [14] reported a study on the deposition chemistry in the HWCVD processes of  $\text{SiH}_4/\text{NH}_3$  mixtures. The densities of  $\text{NH}$ ,  $\text{NH}_2$ ,  $\text{SiH}_3$  radicals as well as  $\text{H}_2$ ,  $\text{N}_2$ , and  $\text{NH}_3$  molecules were measured. It was found that the decomposition efficiency of  $\text{NH}_3$  on the filament surfaces decreases with the introduction of  $\text{SiH}_4$ . Recently, in their study of the plasma chemistry with  $\text{SiH}_4$ - $\text{NH}_3$  system in the process

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of plasma-enhanced CVD, van den Oever et al. [15] have also shown that the densities of N, NH, and NH<sub>2</sub> in the plasma decreased with increasing SiH<sub>4</sub> flow. These studies agree with the fact that a high gas flow ratio of NH<sub>3</sub> to SiH<sub>4</sub>, ranging from 2:1 to greater than 100:1, is needed to get silicon nitride thin films under practical deposition conditions [16, 17]. To understand the interplay of two components in the SiH<sub>4</sub>/NH<sub>3</sub> mixture and the gas-phase chemistry in the HWCVD process using the mixture, we have used laser ionization time-of-flight mass spectrometry (TOF MS) in this work to detect the gas-phase reaction products from a HWCVD reaction chamber using SiH<sub>4</sub>, NH<sub>3</sub>, and SiH<sub>4</sub>/NH<sub>3</sub> mixtures with different pressure ratios.

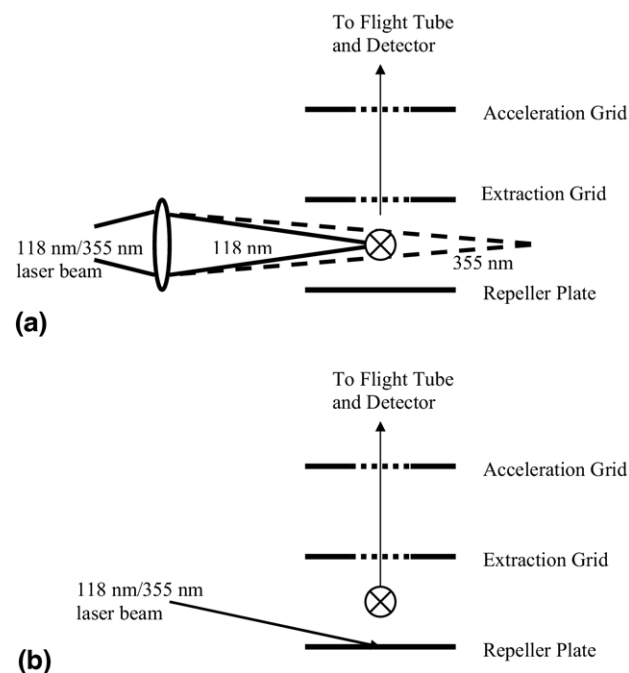
Virtually all species of interest produced in the HWCVD processes using SiH<sub>4</sub> as a source gas have ionization potentials (IP) below 10.5 eV. Therefore, single-photon ionization (SPI) using 118 nm vacuum ultraviolet (VUV) laser radiation has proven to be a powerful technique for examining silane chemistry [6]. Our recent studies have shown that VUV SPI also works well with alkylsilane molecular systems in the HWCVD process [18, 19]. However, many of the products predicted for the NH<sub>3</sub> and SiH<sub>4</sub>/NH<sub>3</sub> systems, notably H<sub>2</sub> and N<sub>2</sub>, have IPs greater than 10.5 eV. A single VUV photon at  $\lambda = 118$  nm cannot ionize them. As a result, a complementary ionization method that would both be capable of ionizing these species and be convenient to implement with a minimum of instrumental modification is highly desirable. Laser-induced electron ionization (LIEI) is an excellent complement to laser ionization [20, 21]. In this method, electron pulses are generated by laser-induced photoemission from a metal target. These photoelectrons are then accelerated in the electric field of the ion source of the TOF MS to ionize the sample molecules. Previous experiments have reported on directing the laser beam to the repeller plate [22], the extraction plate [23], or to a metal target placed in between the two plates [24, 25] in an existing TOF MS. Since the photoelectron energies typically exceed the energy of a single 118 nm photon (10.5 eV), they should be suitable for ionizing all species with IP of either lower or greater than 10.5 eV. In this work, our investigation into the employment of LIEI-MS as an analytical tool in examining gas-phase silicon-nitride film deposition chemistry is reported. The experimental setup for dual ionization source with VUV SPI and LIEI is described. Results from the studies of SiH<sub>4</sub>, NH<sub>3</sub>, and SiH<sub>4</sub>/NH<sub>3</sub> mixtures are presented to illustrate that incorporation of a LIEI source into the conventional laser-ionization system provides a convenient method of simultaneously observing signals of species ionized by both SPI and LIEI.

## Experimental

The 118 nm VUV laser ionization/photoelectron ionization experiments are carried out in a linear TOF mass spectrometer (R. M. Jordan Co., Grass Valley, CA). A

detailed description of the experimental setup using VUV SPI, including the ionization laser source, HWCVD reactor, and the TOF MS, has been reported previously [18, 19, 26]. Briefly, a HWCVD chamber is connected through a 0.15 mm diameter pinhole to the high-vacuum ionization chamber housing the TOF MS. A tungsten filament (0.5 mm diameter, 10 cm length) is placed in the HWCVD reactor and is resistively heated by a DC power supply. The filament temperatures were measured by a two-color pyrometer (Chino Works, Carson, CA). Source gases were introduced into the reactor via a mass flow controller (MKS, type 1179 A) and the pressure in the reactor (7 to 13 torr) was monitored by a capacitance manometer (MKS Baratron, type 626A, Wilmington, MA). The operating pressures in the ionization chamber and flight tube under these reactor conditions were  $4.0 \times 10^{-6}$  torr and  $1.4 \times 10^{-6}$  torr, respectively.

Figure 1 shows the schematic of the dual ionization source used in this work. In SPI mode, chemical species exiting the pinhole of the HWCVD reactor were ionized by a VUV laser beam at 118 nm in a nonresonant process. The 118 nm radiation was produced by frequency tripling the 355 nm UV output from an Nd-YAG laser (Spectra-Physics, LAB-170-10, Mountain View, CA) in a gas cell containing 190 torr of a phase-matched 10:1 Ar/Xe gas mixture. The 355 nm laser energy was kept at 20 mJ/pulse with a repetition rate of 10 Hz. A lithium fluoride (LiF) lens was inserted into the optical path after the gas cell to focus the 118 nm light at the



**Figure 1.** Schematic diagrams of the ionization sources in the ion optics regions of the TOF mass spectrometer for (a) single photon ionization, SPI, by 118 nm VUV photons, and (b) laser-induced electron impact ionization, LIEI, by hitting the repeller plate with the 118 nm/355 nm laser beams.

center of the space between the repeller and extraction grid. The 355 nm UV beam was focused beyond this point. The path of the laser beam was carefully adjusted such that it did not illuminate any metal surfaces in the ion optics.

In LIEI mode, the LiF lens was removed and the path of the laser was adjusted such that the laser beam hit the repeller plate at a shallow angle. Removal of the lens reduced, but did not eliminate, the 118 nm radiation intensity in the ionization region; at the same time it increased the beam divergence. Electrons created directly at the repeller plate should not, however, penetrate very far into the ionization region since they are attracted to the repeller plate which is held at 2.997 kV compared with 2.775 kV to the extraction grid. With the unfocused VUV/UV laser beams, photoelectrons can be produced by straight/scattered/reflected laser light. Being accelerated by the electric field in the ion source, these electrons subsequently gain sufficient energies to ionize the chemical species from the molecular beam exiting the reactor.

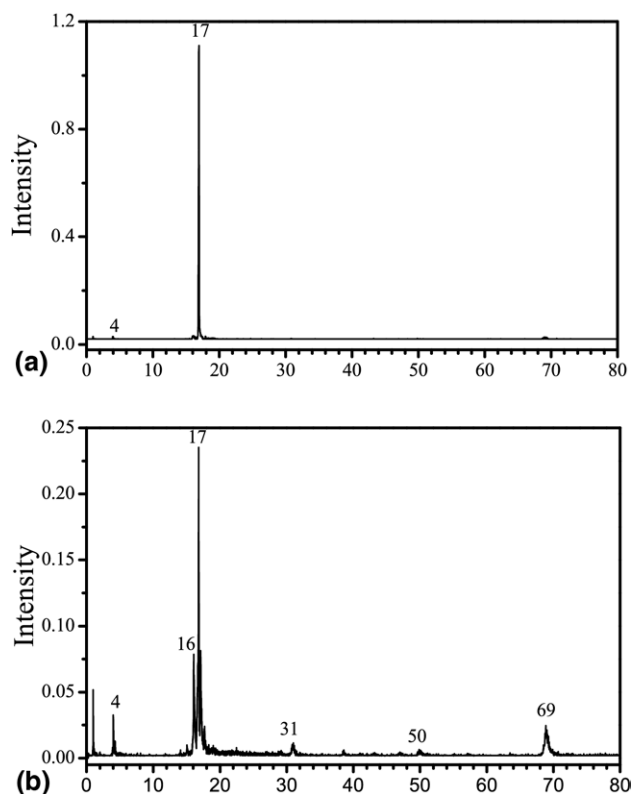
Ions produced by either SPI or LIEI were mass dispersed in a 1 m-long linear TOF mass spectrometer and detected by a microchannel plate (MCP) detector at the end of the flight tube. Signals from the MCP detector were preamplified and displayed on an oscilloscope (Tektronics, TDS 3032B, Markham, Canada). The signals were averaged over 512 laser pulses before being saved to a computer. Pure  $\text{NH}_3$  (99.999%, BOC gases), 1%  $\text{SiH}_4$  in helium (BOC gases), and ultra high purity helium (99.995%, Praxair) were used to prepare the  $\text{SiH}_4/\text{NH}_3/\text{He}$  mixtures in a 2.25 L sample cylinder.

## Results and Discussions

Source gaseous mixtures, including 50%  $\text{NH}_3$  in helium (He), 1%  $\text{SiH}_4$  in He, and  $\text{NH}_3/\text{SiH}_4$  mixtures in He with different pressure ratios of  $\text{NH}_3$  to  $\text{SiH}_4$ , were studied in this work. The mass spectra from  $\text{NH}_3$  and  $\text{SiH}_4$  gases before and after the filament was turned on using both SPI and LIEI methods were recorded to evaluate the effectiveness of our LIEI source.

### $\text{NH}_3$ and $\text{SiH}_4$ Components

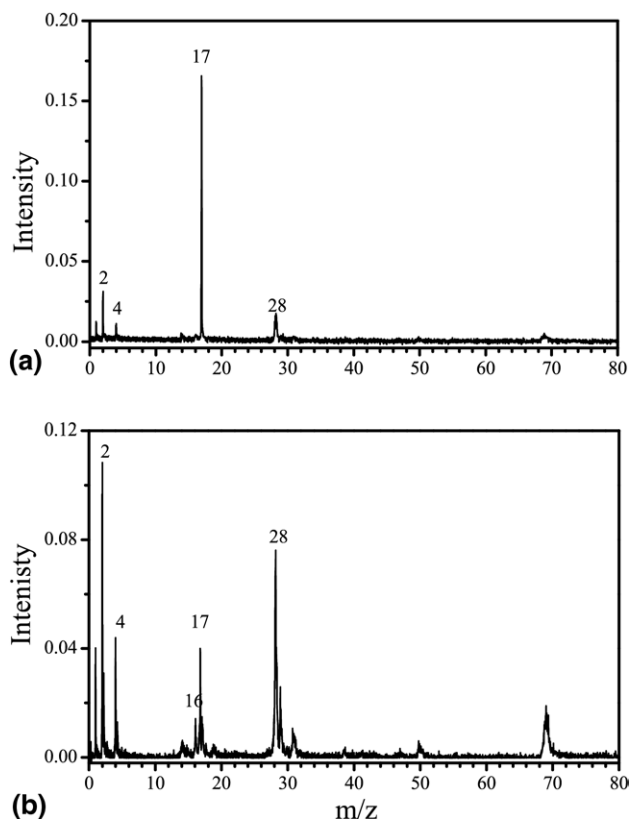
The room-temperature mass spectra of 50%  $\text{NH}_3/\text{He}$  recorded under SPI and LIEI conditions are shown in Figure 2a and b, respectively. The mass spectrum recorded using SPI is dominated by the parent mass at  $m/z = 17$ . All other peaks including the  $\text{He}^+$ ,  $\text{NH}_2^+$ , and  $\text{NH}^+$  are very weak. These weak peaks are attributed to electron impact (EI) ionization due to a small amount of photoelectrons caused by scattered UV radiation in the photoionization region. By changing to the LIEI method, the  $\text{He}^+$ ,  $\text{NH}_2^+$ , and  $\text{NH}^+$  peaks are greatly enhanced, as expected. One more fragment peak of  $\text{N}^+$  was also observed. The IP of He is 24.6 eV. The appearance potentials (AP) of  $\text{NH}_2^+$ ,  $\text{NH}^+$ , and  $\text{N}^+$  from  $\text{NH}_3$  are 15.76 eV, 17.1 eV, and 22.6 eV, respec-



**Figure 2.** Room temperature mass spectra of 50%  $\text{NH}_3/\text{He}$  recorded using (a) SPI mode; and (b) LIEI mode.

tively. All this indicates that EI source becomes important in the LIEI mode where the LiF lens is pulled out of the optical path and the laser beam is guided toward the repeller plate. Also of interest is the fact that the baseline is much less noisy when the SPI setup is used, a fact that has been observed by others [27]. When the LIEI source is used, the baseline is much noisier and background peaks at  $m/z = 1$ , 31, 50, and 69 appear with increased intensity. It is also noted from Figure 2 that by changing from the SPI to LIEI source, the intensity of the  $\text{NH}_3$  peak is significantly reduced. The reason will be discussed later.

Figure 3 shows the comparison of the TOF mass spectra recorded for 12 torr of 50%  $\text{NH}_3/\text{He}$  mixture in the HWCVD reactor after the filament was turned on for 2 min at 1500 °C under SPI and LIEI conditions. Compared with the room-temperature mass spectra of the same mixture, clearly, after the filament was turned on, the intensity of  $\text{NH}_3$  peak decreases, while at the same time two new peaks at  $m/z = 2$  and 28 increase in intensity. These two new peaks correspond to  $\text{H}_2$  and  $\text{N}_2$  molecules, respectively, which are produced from hot-wire decomposition of  $\text{NH}_3$  in the reactor. In the work by Umemoto et al. [12], it was found that the primary decomposition products of  $\text{NH}_3$  on the W filament are  $\text{NH}_2$  and H radicals. They also detected the stable  $\text{N}_2$  and  $\text{H}_2$  molecules using mass spectrometer. Melton and Emmett [28] show that the N-containing species in the catalytic decomposition of  $\text{NH}_3$  over Pt at



**Figure 3.** Mass spectra of 50% NH<sub>3</sub>/He sample after the W filament was turned on for 2 min at 1500 °C recorded with (a) SPI mode; and (b) LIEI mode.

a pressure of 0.1 torr and a temperature of 1200 °C are N<sub>2</sub>, NH<sub>2</sub>, NH, and N, with N<sub>2</sub> as the dominant of all. NH<sub>2</sub> is the most abundant radical species detected among the three. One of the important NH<sub>2</sub> loss channel proposed by Stohard et al. [29] is the heterogeneous reactions on the chamber wall, which is found to be much faster than the gas-phase biradical combination reactions. In our experiment, the dominant species when the W filament is on are the stable H<sub>2</sub> and N<sub>2</sub> molecules. They are believed to be formed by reactions on the chamber walls [26].

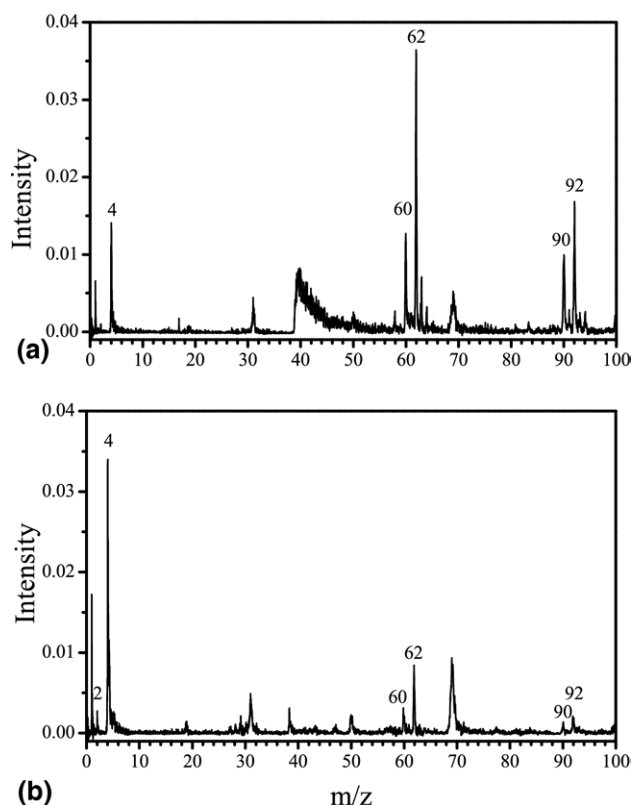
The impact of the LIEI source on the observed mass spectra is dramatic as seen from comparing Figure 3a and b. In using LIEI, the intensities of the two product peaks, i.e., H<sub>2</sub> and N<sub>2</sub>, are significantly enhanced. This is clearly more representative of the product distributions in the reactor with NH<sub>3</sub> gas in. The mass resolution ( $m/\Delta m$ ) of the H<sub>2</sub>, N<sub>2</sub>, and He peaks observed with LIEI source is found to be not higher than 115, which is much less than the mass resolution for the NH<sub>3</sub> peak recorded using SPI method, typically of 200. This implies that the enhanced signals of H<sub>2</sub>, N<sub>2</sub>, and He come from the electron impact ionization source. As mentioned before, the NH<sub>2</sub><sup>+</sup> ( $m/z$  = 16) and NH<sup>+</sup> ( $m/z$  = 15) peaks were observed in the room-temperature mass spectra of NH<sub>3</sub>. The intensities of these two peaks decreased with their parent ion, NH<sub>3</sub><sup>+</sup>, after the fila-

ment was turned on. Examination of the intensity ratios of NH<sub>2</sub><sup>+</sup>/NH<sub>3</sub><sup>+</sup> and NH<sup>+</sup>/NH<sub>3</sub><sup>+</sup> from the entire filament on mass spectra of NH<sub>3</sub> shows no obvious increase. Therefore, the NH<sub>2</sub><sup>+</sup> and NH<sup>+</sup> signals should originate from the dissociation of NH<sub>3</sub> upon ionization, and they are not related to dissociation of NH<sub>3</sub> on the hot W filament. This is not surprising since the NH<sub>2</sub> and NH radicals produced directly from the hot filament cannot survive the sampling conditions in our HWCVD reactor. However, the intensity of the peak at  $m/z$  = 14 is seen to increase after the filament is on. The peak is most likely the N<sup>+</sup> signal from the dissociation of N<sub>2</sub><sup>+</sup> upon ionization. This is supported by the fact that the time distribution curves of two peaks at  $m/z$  = 28 and 14 are very similar.

At a filament temperature of 1500 °C, the NH<sub>3</sub> peak intensity reached a steady-state after about 3 min. The intensities of the two products, N<sub>2</sub> and H<sub>2</sub>, also reached their respective steady-state at the same time. With increasing filament temperatures, the time to reach the steady-state intensities is shortened. At  $T \geq 1800$  °C, the steady-state intensity is achieved right after the filament is on. At the steady state, the intensity ratio of the three peaks is found to be 0.0207:0.715:1.00 for NH<sub>3</sub>:N<sub>2</sub>:H<sub>2</sub>. By using the reference mass spectra of a mixture sample with H<sub>2</sub>, N<sub>2</sub>, and NH<sub>3</sub> in equal quantities recorded under the same LIEI conditions, the true ratio of the three molecules in the HWCVD reactor can be determined to be 0.0130:0.291:1.00. This suggests that at the steady-state in the HWCVD reactor, the amount of H<sub>2</sub> is the greatest among the three stable molecules, the N<sub>2</sub> steady-state intensity is 29.1% that of H<sub>2</sub>, and NH<sub>3</sub> is only 1.30% that of H<sub>2</sub>. The average decomposition percentage of NH<sub>3</sub> when using the 50% NH<sub>3</sub>/He sample is found to be 98.8%, close to a complete decomposition.

For the SiH<sub>4</sub> component, the principal gas-phase reaction products were found to be H<sub>2</sub>, Si<sub>2</sub>H<sub>6</sub> ( $m/z$  = 62), and Si<sub>3</sub>H<sub>8</sub> ( $m/z$  = 92) from our experiments with 12 torr of 1% SiH<sub>4</sub>/He mixture. This is shown in Figure 4, the mass spectra of 1% SiH<sub>4</sub> recorded right after the filament was turned on at 1500 °C. The two peaks observed at  $m/z$  = 60 and 90 are believed to be formed from photoionization of Si<sub>3</sub>H<sub>8</sub> and Si<sub>4</sub>H<sub>10</sub>, respectively, formed in the reactor [9]. All other mass peaks except that at  $m/z$  = 4 (He<sup>+</sup>) are from background. The broad peak at an  $m/z$  value of around 40 is found to be due to the filament on process. The reason for the observation of this broad feature is currently unclear. It can be seen from Figure 4 that the intensities of the Si<sub>2</sub>H<sub>6</sub><sup>+</sup> (IP = 9.7 eV) and Si<sub>3</sub>H<sub>8</sub><sup>+</sup> (IP = 9.2 eV) are enhanced when the SPI method was used. Similar to what was observed for 50% NH<sub>3</sub>, the H<sub>2</sub><sup>+</sup> and He<sup>+</sup> peaks were stronger when using the LIEI method. It is noted that a weak photo-fragment peak of SiH<sub>2</sub><sup>+</sup> ( $m/z$  = 30) from SiH<sub>4</sub> (AP = 11.6 eV) were observed using either SPI or LIEI, and its intensity was found to be enhanced when using LIEI mode. After the filament was turned on, the SiH<sub>2</sub><sup>+</sup> peak intensity decayed rapidly with filament on time for each temperature tested. The same intensity was also





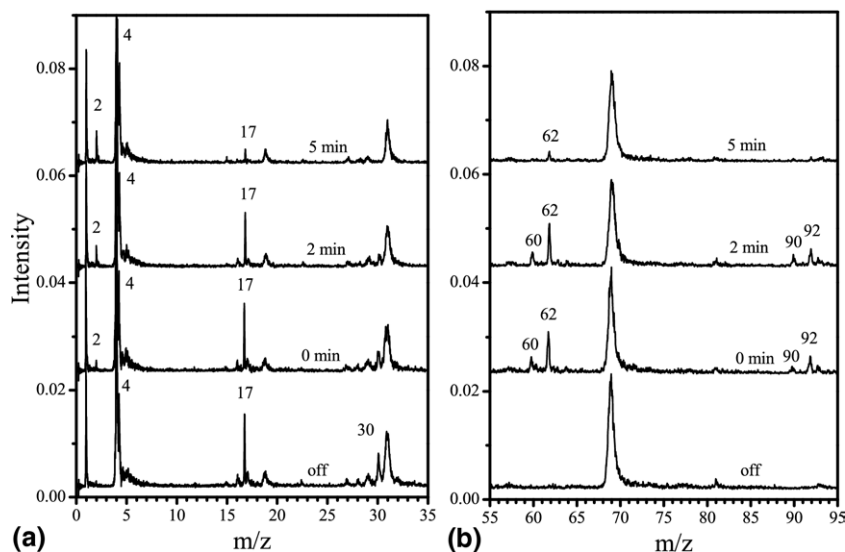
**Figure 4.** Mass spectra of 1%  $\text{SiH}_4/\text{He}$  sample right after the W filament was turned on ( $t = 0$  min) at  $1500^\circ\text{C}$  recorded with (a) SPI mode; and (b) LIEI mode.

found to decrease with increasing filament temperatures. Therefore, the enhanced  $\text{SiH}_2^+$  ( $m/z = 30$ ) signal observed with LIEI source provides good indication of the  $\text{SiH}_4$  depletion in the reactor since its parent ion is difficult to observe.

In SPI mode, the 118 nm photons are focused directly into the molecular beam, and so the intensity of 10.5 eV radiation is high. Under these conditions, it would be expected that any peak corresponding to a species that can be ionized by the 118 nm photons ( $\text{IP} < 10.5$  eV) would increase in intensity. The corollary of this is that any species that cannot be ionized by 10.5 eV radiation would experience a decrease in peak intensity. Examination of the behavior of the  $\text{SiH}_4\text{--NH}_3$  samples demonstrates that this is indeed the case. The best illustration of the effect can be made using the  $\text{NH}_3$ ,  $\text{Si}_2\text{H}_6$ ,  $\text{H}_2$ , and  $\text{N}_2$  peaks. Both  $\text{NH}_3$  and  $\text{Si}_2\text{H}_6$  have IPs below 10.5 eV, and both show a greater intensity in SPI mode compared with LIEI mode.  $\text{H}_2$  and  $\text{N}_2$ , by contrast, are significantly weaker when SPI mode is employed. Evidently any species capable of being ionized by 10.5 eV radiation has a higher intensity in SPI mode, whereas species with  $\text{IP} > 10.5$  eV are only effectively ionized when LIEI is used.

### $\text{SiH}_4/\text{NH}_3$ Mixtures

The experiments on the two pure components, i.e.,  $\text{SiH}_4$  and  $\text{NH}_3$ , have shown that the LIEI source coupled with TOF MS provides a useful tool to diagnose the gas-phase chemical species involved with both  $\text{SiH}_4$  and  $\text{NH}_3$  in the HWCVD reactor. Therefore, it is the method of choice to study the behavior of  $\text{SiH}_4/\text{NH}_3$  mixture in the same reactor. Figure 5 shows the mass spectra recorded for the 1:1  $\text{NH}_3\text{:SiH}_4$  mixture after the filament was turned on at  $1500^\circ\text{C}$  for 0, 2, and 5 min, respectively. The room-temperature spectrum of the same mixture is shown for comparison. The peaks representative of the  $\text{SiH}_4$  chemistry in the reactor, i.e., the  $\text{Si}_2\text{H}_6$  ( $m/z = 62$ ),  $\text{Si}_3\text{H}_8$  ( $m/z = 92$ ), and  $\text{H}_2$ , were observed from the mass spectra. Similar to 1%  $\text{SiH}_4/\text{He}$  sample,



**Figure 5.** The LIEI mass spectra of the  $\text{NH}_3/\text{SiH}_4$  mixture [ $\text{P}(\text{NH}_3)\text{:P}(\text{SiH}_4) = 1\text{:}1$ ] recorded 0, 2, and 5 min after the filament was on at  $1500^\circ\text{C}$  in mass regions (a) between 0 and 35 amu; and (b) between 55 and 95 u.

the peaks at  $m/z = 62$  and  $92$  decay with increasing filament-on time, whereas the  $H_2$  peak increases to a steady-state intensity. On the other hand, the characteristic product peak of  $N_2$  from  $NH_3$  chemistry was not observed for this mixture even after the filament was turned on at  $1500\text{ }^\circ\text{C}$  for 5 min. At this temperature, for the pure  $NH_3$  component, the  $N_2$  steady-state intensity was achieved 3 min after the filament was turned on. Therefore, it can be concluded that when equal quantity of  $SiH_4$  and  $NH_3$  is mixed,  $SiH_4$  chemistry is dominant and the  $NH_3$  chemistry is suppressed by the addition of  $SiH_4$ . At a raised filament temperature of  $1800\text{ }^\circ\text{C}$ , signals from  $N_2$  peak start to come out, indicating a lesser extent of suppression at higher temperatures. Although the  $N_2$  peak was not observed at  $T < 1800\text{ }^\circ\text{C}$ , the  $NH_3$  peak intensity was found to decrease. For example, right after the filament was turned on at  $1500\text{ }^\circ\text{C}$  and  $1800\text{ }^\circ\text{C}$ , the  $NH_3$  peak intensity is decreased by 81.2% and 87.5%, respectively. It is known that gas-phase reactions, such as  $H + NH_3$  and  $SiH_4 + NH_3$ , are slow [30, 31]; therefore, the consumption of  $NH_3$  in the secondary gas-phase reactions is less likely. The decrease in the  $NH_3$  signal is probably due to surface reactions.

Various other  $SiH_4/NH_3$  mixtures with increasing  $NH_3$  contents were studied in this work. The compositions of these mixtures are listed in Table 1. At the lowest filament temperature of  $1500\text{ }^\circ\text{C}$  tested, no mixtures show a decrease in the  $NH_3$  peak intensity for the full 10-min run except that the 200:1  $NH_3:SiH_4$  mixture shows an 11.3% decrease at a filament-on time of 10 min compared with its room-temperature intensity. With increasing filament temperature and filament-on time, the  $NH_3$  peak intensity starts to decrease. Table 2 lists the  $NH_3$  decomposition percentage for the 200:1, 100:1, 50:1, and 20:1  $NH_3/SiH_4$  mixtures 0, 5, and 10 min, respectively, after the filament is on at  $1500\text{ }^\circ\text{C}$ ,  $1700\text{ }^\circ\text{C}$ , and  $2000\text{ }^\circ\text{C}$ . The experimental error for the  $NH_3$  decomposition percentages was determined to be no greater than 4.9 from the measurements of 50%  $NH_3/He$  sample. Comparison of the 50%  $NH_3/He$  and 100:1  $NH_3/SiH_4/He$  samples clearly shows the negative effect of  $SiH_4$  addition on the decomposition of  $NH_3$ . Only at a filament temperature of  $2000\text{ }^\circ\text{C}$  do the two samples show comparable behavior regarding the  $NH_3$  decomposition. By comparing the 200:1, 50:1, and

**Table 2.** The decomposition percentage of  $NH_3$  for the four  $NH_3/SiH_4$  mixtures and 50%  $NH_3/He$

	50% $NH_3$	200: 1	100:1	50: 1	20:1
1500 $^\circ\text{C}$ at 0 min	37.1	0	0	0	0
1500 $^\circ\text{C}$ at 5 min	98.0	0	0	0	0
1500 $^\circ\text{C}$ at 10 min	99.4	11.0	0	0	0
1700 $^\circ\text{C}$ at 0 min	79.4	0	0	0	0
1700 $^\circ\text{C}$ at 5 min	98.6	76.1	29.2	32.2	0
1700 $^\circ\text{C}$ at 10 min	99.2	96.1	54.7	62.1	4.9
2000 $^\circ\text{C}$ at 0 min	92.8	62.0	29.2	57.4	19.4
2000 $^\circ\text{C}$ at 5 min	99.1	98.6	99.0	98.9	98.9
2000 $^\circ\text{C}$ at 10 min	99.2	99.0	99.1	99.0	99.4

20:1  $NH_3/SiH_4/He$  samples, it can be seen that the suppression of  $SiH_4$  on the  $NH_3$  decomposition is enhanced with the increasing amount of  $SiH_4$  in the mixture.

For all the mixtures listed in Table 1, when the  $NH_3$  peak shows a decrease in intensity, the  $N_2$  peak is observed in the mass spectra. Below filament temperatures of  $1700\text{ }^\circ\text{C}$ , if observed, the intensity of the  $N_2$  peak from all the mixture keeps increasing with filament-on time. No steady-state intensity is reached at the end of the 10-min run for each tested temperature lower than  $1700\text{ }^\circ\text{C}$ . This is in sharp contrast with the behavior of the pure  $NH_3$  component. At temperatures higher than  $1800\text{ }^\circ\text{C}$ , the  $N_2$  intensity reaches a steady-state for all mixtures. However, the time it takes to reach the steady-state intensity is different, as shown in Table 1. With more  $SiH_4$  in the mixture, the time taken to achieve the  $N_2$  steady-state intensity is increased. This supports the conclusion that more  $SiH_4$  molecules suppress the  $NH_3$  decomposition to an even greater extent.

Previous study on the deposition chemistry of  $SiH_4/NH_3$  mixtures in HWCVD [14] also showed the negative effect of  $SiH_4$  addition on the  $NH_3$  decomposition. The reason was ascribed to the competition of active catalytic sites on the filament between  $SiH_4$  and  $NH_3$ . In our experiments with the  $SiH_4/NH_3$  mixtures, consistent observation was made that an increased current was needed to maintain a constant filament temperature. At certain point, a temperature jump occurred, accompanied by the rapid decrease of the  $NH_3$  peak intensity and strong build-up of the  $N_2$  and  $H_2$  peak intensities in the mass spectra. This suggests that  $SiH_4$  interacts with the filament more aggressively than  $NH_3$ , forming silicide on the filament, and this is why more power is needed to maintain the filament temperature. The presence of silicon element on the surface of W filament exposed to the  $SiH_4/NH_3$  mixtures is confirmed using the electron microprobe analysis (EMPA) (JEOL JXA-8200) of filament. When  $NH_3$  takes over the filament at a later time and with the power build-up on the filament, the electrical properties and the emissivity of the filament changed, causing a temperature spike. The access of  $NH_3$  to the filament decomposes the molecule and the corresponding stable

**Table 1.** The composition of various  $SiH_4/NH_3$  mixtures and the filament-on time to reach the  $N_2$  steady-state intensity

Mixtures	$P(NH_3)^a$ (torr)	$P(SiH_4)$ (torr)	Time (min)		
			1800 $^\circ\text{C}$	1900 $^\circ\text{C}$	2000 $^\circ\text{C}$
50% $NH_3$	1000	0	0	0	0
200: 1	660	3.3	3	2	1
100: 1	1000	10	3	3	2
50: 1	660	13.5	3	3	2
20: 1	380	16.3	5	3	2

<sup>a</sup>All mixtures are balanced with helium to a total pressure of 3 atm.

products,  $H_2$  and  $N_2$ , are then detected by our mass spectrometer. Therefore, our experimental observations also suggest that when  $SiH_4$  and  $NH_3$  are mixed in the HWCVD reactor, they tend to compete for the active site on the filament.  $SiH_4$  is clearly more dominant in this regard.

### *The Ionization Process*

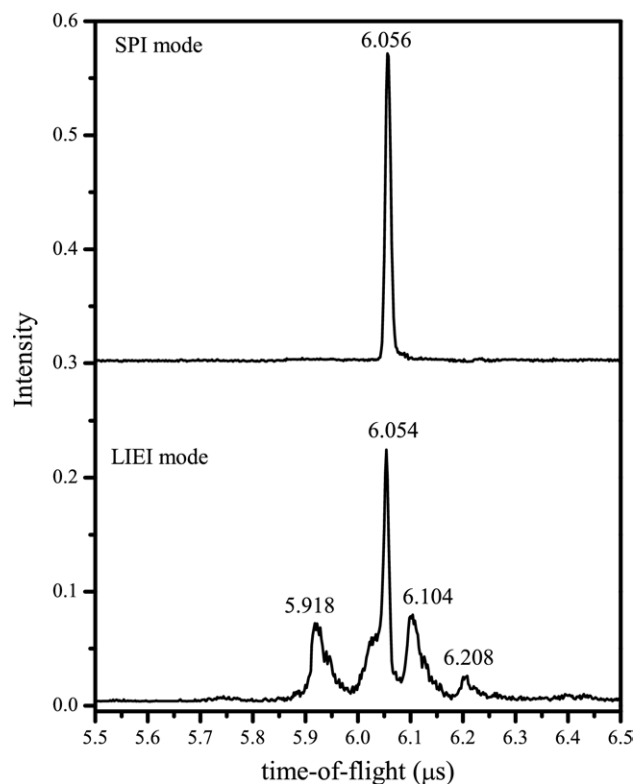
The capability of the LIEI source for detecting the gas-phase chemical products formed in the HWCVD processes with  $SiH_4$ ,  $NH_3$  and their mixtures has been clearly demonstrated by the experimental results presented above. Effort is also made in our study to understand the nature of the ionization process. In the LIEI mode, the 355 nm UV beam passes the frequency-tripling gas cell filled with Ar/Xe before it hits the repeller plate. The two wavelengths, 355 nm and 118 nm, are not separated before the beam's arrival in the ionization chamber. Although the 118 nm VUV photons have much higher energy than the 355 nm UV photons, the intensity of VUV light is much lower due to the low conversion efficiency of about  $10^{-5}$  [32]. In our experiments, it was found that signals from  $Na^+$  and  $K^+$  were observed in a few instances or when the 355 nm UV light intensity was increased; however, other metal ion peaks such as the iron ion peak were absent. These alkali metals have work functions below 3.5 eV, whereas the work functions of the main components in stainless steel, e.g., iron and nickel, are above 3.5 eV. This suggests that the 355 nm radiation is a much more likely candidate responsible for photoelectron emission in the LIEI mode. It is noted from Figure 2a and Figure 3a that in the SPI mode, where care was taken to prevent the laser light from hitting any surfaces, signals originating from LIEI, e.g.,  $H_2^+$ ,  $He^+$ , and  $N_2^+$ , are still observable, but at reduced intensities. In this mode, when the tripling medium, Xe gas, was pumped out of the gas cell, the  $NH_3$  peak in Figure 2a and Figure 3a is completely gone, as expected. However, the intensities of the  $H_2^+$ ,  $He^+$ , and  $N_2^+$  peaks were found to stay the same as those with the Xe gas in the cell. This clearly demonstrates that the weak LIEI signals present in the mass spectra recorded using the SPI mode is caused by scattered 355 nm UV light in the ionization region. Therefore, it is believed that the 355 nm UV beam is responsible for producing pulsed electrons beam in both ionization methods.

In the LIEI mode, although the 355 nm UV laser beam was directed to the repeller plate held at much higher positive potential than the extraction grid in the TOF MS, the results presented above have shown that the LIEI source works well. Previous work by Schriemer and Li [22] has also illustrated that ion signals were observed by striking the repeller plate with a laser beam. It is unlikely that the repeller would be the main source of ionizing electrons. The small kinetic energy the electrons obtain in the laser-induced photoelectron emission process is not enough to even drive them 0.1

mm into the ionization region and it is impossible for them to be accelerated to energies above 10.5 eV. More probable is that the electrons originate from straight/reflected/scattered laser light hitting the extraction grid and/or other metal surfaces as seen by Colby and Reilly [20]. Such electrons can then be accelerated to hundreds of eV by the extraction field between the repeller plate and extraction grid. Under the LIEI mode used in our experiment, the LiF lens was pulled out of the optical path, therefore, the laser light is divergent. A measurement of the UV beam divergence under the condition of LIEI mode indicates a beam spot size with a diameter of about 15 mm, which is larger than the spacing of 12.7 mm between the repeller and the extraction plates. Therefore, the divergent UV beam can hit the extraction plate to produce photoelectrons that can be used for molecular ionization. In our current experimental setup, it is not likely to guide the 355 nm UV beam toward the extraction plate directly and still be able to generate the 118 nm VUV laser radiation. A more controlled generation of photoelectrons will be necessary before a more complete understanding of the ionization process can be obtained. Such experiments are currently underway.

Comparison between the SPI and the LIEI source indicates that SPI tends to give mass peaks with a good resolution while the peaks from the LIEI source are broader. This is probably due to the broad initial spatial distribution of the ions generated in the LIEI source, as has been discussed in previous work [23, 25]. In the LIEI mode used in our experiment, two laser beams with different wavelengths are present in the ionization region; therefore, the species with IPs below 10.5 eV can be ionized simultaneously by SPI and LIEI sources. For the  $NH_3$  peak (IP: 10.07 eV), it is interesting to note that the peaks due to SPI and LIEI for the same mass were somewhat separated in time as can be seen in Figure 6. The  $NH_3$  ions due to LIEI ( $t = 6.104 \mu s$ ) arrive at the detector a bit later than those due to SPI ( $t = 6.054 \mu s$ ). The peak at  $t = 5.918 \mu s$  is from  $NH_2^+$  and the source for the peak at  $t = 6.208 \mu s$  is unknown. The average intensity ratio of the  $NH_2^+$ ,  $NH^+$  peak, respectively, to the  $NH_3^+$  due to LIEI is found to be 0.89:1 and 0.088:1. This is comparable to the ratio of 0.81:1 and 0.076:1 for  $NH_2^+/NH_3^+$  and  $NH^+/NH_3^+$  from the conventional 70 eV EI MS of  $NH_3$ , implying the electron energy in the LIEI source should be close to 70 eV. Considering the electric field of 175 V/cm between the repeller plate and the extraction grid, the ionization would occur in a spatial region closer to the extraction plate, consistent with the origin of the photoelectrons being from the extraction plate. This may also explain why the  $NH_3$  ions due to the energetic electrons arrive at a later time than the same ions from VUV photoionization, which occurs at the center of the repeller plate and the extraction grid.

A characteristic feature of the current LIEI source is the presence of both 118 nm VUV and 355 nm UV radiations in the ionization region, leading to the coex-



**Figure 6.** Room temperature mass spectra of the  $\text{NH}_3/\text{SiH}_4$  mixture [ $\text{P}(\text{NH}_3):\text{P}(\text{SiH}_4) = 100:1$ ] in the time-of-flight region of  $\text{NH}_3$  recorded under SPI and LIEI modes.

istence of SPI and LIEI mode. A valuable consequence is that while it enables us to observe the species with IPs above 10.5 eV, the peaks with IPs below 10.5 eV are visible with great intensity. As a result, the current LIEI source enables us to simultaneously observe all species of interest. As mentioned previously, SPI is a “soft” ionization method having a great advantage of generating molecular ions with relatively little fragmentation of the parent molecule. The EI source complements the SPI source in the sense that a “fingerprint” fragmentation pattern can be obtained for structure identification. An ionization source allowing the simultaneous observation of the molecular ion with great intensity (from SPI) and its fragment ions (from EI) would be advantageous, especially for monitoring reactions in which numerous species may appear at a given  $m/z$  value. In this way, an even more “universal” technique than 118 nm SPI alone is developed, and it is possible to identify unknowns using reference electron impact spectra from databases and also its parent mass. For the  $\text{SiH}_4\text{--NH}_3$  system, for example, the observation of  $\text{NH}_4^+$ ,  $\text{NH}_2^+$  signals together with an enhanced  $\text{NH}_3^+$  parent ion provide additional information about the parent ion. The increased intensity in the peak at  $m/z = 14$  representing  $\text{N}^+$  with the product peak at  $m/z = 28$  confirms the assignment of the peak to  $\text{N}_2$ . It is anticipated that with increasing complexity of the source molecular system, the observation of both parent ions and its

fragmentation peaks should provide significant help in positive identification of structure isomers produced in the complex HWCVD systems.

## Conclusions

The gas-phase chemical species produced from a HWCVD reactor with  $\text{SiH}_4$ ,  $\text{NH}_3$ , and their mixtures with different pressure ratios were examined using the conventional 118 nm VUV SPI source and a newly-developed LIEI source. Comparison of both ionization sources has shown that species with an IP below 10.5 eV, such as  $\text{Si}_2\text{H}_6$ ,  $\text{Si}_3\text{H}_8$ , and  $\text{NH}_3$ , exhibited stronger peaks in the mass spectra recorded under SPI conditions. Species with IPs above 10.5 eV, such as  $\text{H}_2$ ,  $\text{N}_2$  and He, were only weakly visible under the same conditions. The incorporation of the LIEI source significantly enhanced the intensity of those species with IPs above 10.5 eV. Consequently, the mass spectra recorded with LIEI source is more representative of the product distributions for the source gas systems under study. A unique feature of the developed LIEI source is the coexistence of the SPI source due to the 118 nm VUV light. In using this ionization source, a simultaneous observation of parent ions with enhanced intensity from VUV SPI and their “fingerprint” fragmentation ions from LIEI has been achieved. This should provide a useful ionization source for distinguishing structure isomers often present in a complex HWCVD system. More importantly, this is achieved in one ionization source; no switch between SPI and LIEI modes via apparatus modification is needed.

The LIEI source developed proves to be very useful in studying the gas-phase chemistry in the process of HWCVD with the  $\text{SiH}_4\text{--NH}_3$  system. Examination of a  $\text{SiH}_4/\text{NH}_3$  mixture with a pressure ratio of 1:1 shows that  $\text{SiH}_4$  chemistry is predominant when equal quantity of  $\text{SiH}_4$  and  $\text{NH}_3$  is mixed. The  $\text{N}_2$  product characteristic of  $\text{NH}_3$  chemistry in the reactor is absent with this mixture at low temperatures, indicating a negative effect from the  $\text{SiH}_4$  addition on the  $\text{NH}_3$  chemistry. Study of various other  $\text{SiH}_4/\text{NH}_3$  mixtures with a pressure ratio ( $\text{P}_{\text{NH}_3}:\text{P}_{\text{SiH}_4}$ ) of 20:1, 50:1, 100:1, and 200:1 confirms that the  $\text{NH}_3$  decomposition is suppressed by the presence of  $\text{SiH}_4$  molecule in the mixture. This is in agreement with previous studies and explains why a high flow ratio of  $\text{NH}_3$  to  $\text{SiH}_4$  is needed under practical deposition conditions to produce silicon nitride thin films. The  $\text{NH}_3$  decomposition percentage is found to decrease and the time for  $\text{N}_2$  and  $\text{H}_2$  to reach their steady-state intensities to increase with increasing  $\text{SiH}_4$  content in the mixture. This suggests that the extent to which the  $\text{NH}_3$  decomposition is suppressed is enhanced with more  $\text{SiH}_4$  in the system. Currently in our laboratory, experiments for a more controlled photoelectron generation is underway for a clear understanding of the LIEI process involved. In addition, our preliminary study has indicated that several possible structures could be responsible for the same mass peak



observed when trimethylsilane is used as a source gas in a HWCVD reactor. The application of the LIEI source to this more complicated source gas system for positive identification of structure isomers will be explored in the future.

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